Rainbows, Quantum Billiards, and the Birth of Reflections Stokes Phenomenon Exemplified

Will Hoffer

University of California, Riverside

math@willhoffer.com

October 22, 2020

▲□▶ ▲□▶ ▲□▶ ▲□▶ □ のQ@



• What is the Stokes Phenomenon?



- What is the Stokes Phenomenon?
- Example One: Rainbows/The Airy function

- What is the Stokes Phenomenon?
- Example One: Rainbows/The Airy function
- Example Two: Birth of Reflections/Helmholtz Equation

- What is the Stokes Phenomenon?
- Example One: Rainbows/The Airy function
- Example Two: Birth of Reflections/Helmholtz Equation
- Example Three: Quantum Billiards & Weyl Expansions

・ロト ・ 日 ・ ・ 日 ・ ・ 日 ・ ・ りへぐ

Asymptotic Expansions, in the sense of Poincaré The following are equivalent. As $z \to \infty$:

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 のへぐ

$$f(z) \sim \sum_{n=1}^{\infty} a_n \frac{1}{z^n}$$

Asymptotic Expansions, in the sense of Poincaré The following are equivalent. As $z \to \infty$:

$$\begin{split} f(z) &\sim \sum_{n=1}^{\infty} a_n \frac{1}{z^n} \\ f(z) &= \sum_{n=1}^{N} a_n \frac{1}{z^n} + O\left(\frac{1}{z^{N+1}}\right) \end{split}$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 のへぐ

Asymptotic Expansions, in the sense of Poincaré The following are equivalent. As $z \to \infty$:

$$f(z) \sim \sum_{n=1}^{\infty} a_n \frac{1}{z^n}$$
$$f(z) = \sum_{n=1}^{N} a_n \frac{1}{z^n} + O\left(\frac{1}{z^{N+1}}\right)$$
$$f(z) = \sum_{n=1}^{N} a_n \frac{1}{z^n} + o\left(\frac{1}{z^N}\right)$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

Let $\sinh(z) = \frac{1}{2}(e^z - e^{-z})$ where $z \in \mathbb{C}$.



Let $\sinh(z) = \frac{1}{2}(e^z - e^{-z})$ where $z \in \mathbb{C}$.



◆□▶ ◆□▶ ◆三▶ ◆三▶ ○□ ● のへで

Let $\sinh(z) = \frac{1}{2}(e^z - e^{-z})$ where $z \in \mathbb{C}$. Observe, as $z \to \infty$: $\sinh(z) \sim \begin{cases} \frac{1}{2}e^z & \Re(z) > 0\\ -\frac{1}{2}e^{-z} & \Re(z) < 0 \end{cases}$

▲□▶ ▲□▶ ▲□▶ ▲□▶ □ のQ@



Let $\sinh(z) = \frac{1}{2}(e^z - e^{-z})$ where $z \in \mathbb{C}$. Observe, as $z \to \infty$: $\sinh(z) \sim \begin{cases} \frac{1}{2}e^z & \Re(z) > 0\\ -\frac{1}{2}e^{-z} & \Re(z) < 0 \end{cases}$





▲□▶ ▲□▶ ▲□▶ ▲□▶ □ のQ@

Observe the change in behavior across the imaginary axis.

$$f(z) = \sinh(z)$$



◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 のへぐ

Observe the change in behavior across the imaginary axis.



Observe the change in behavior across the imaginary axis.



◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

Broadly speaking, the Stokes phenomenon is that asymptotic expansions may change behavior in the complex plane.

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 のへぐ

Broadly speaking, the Stokes phenomenon is that asymptotic expansions may change behavior in the complex plane.

More strictly, a Stokes phenomenon is a change arising from the "conception and subsequent birth" of terms that appear and become active with the changing phase.

▲□▶ ▲□▶ ▲□▶ ▲□▶ □ のQ@

Broadly speaking, the Stokes phenomenon is that asymptotic expansions may change behavior in the complex plane.

More strictly, a Stokes phenomenon is a change arising from the "conception and subsequent birth" of terms that appear and become active with the changing phase.

R. B. Dingle's Description

At a certain phase drawn in the complex plane as a "Stokes ray", an "associated function" appears, disappears or changes its numerical multiplier."

うして ふゆ く は く は く む く し く

Supernumerary Rainbow



◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 - のへで

The Airy Model



・ロト ・ 一下・ ・ ヨト・・

= 900

The Airy Model





Here, U denotes a solution to the Airy equation, $\frac{d^2U}{dn^2}+\frac{n}{3}U=0.$

$$U = An^{-\frac{1}{4}} e^{\frac{2}{3}\sqrt{-\frac{n^2}{3}}} \left\{ 1 - \frac{1 \cdot 5}{1} \frac{\sqrt{-1}}{16\sqrt{(3n^3)}} + \frac{1 \cdot 5 \cdot 7 \cdot 11}{1 \cdot 2} \left(\frac{\sqrt{-1}}{16\sqrt{(3n^3)}} \right)^2 - \frac{1 \cdot 5 \cdot 7 \cdot 11 \cdot 13 \cdot 17}{1 \cdot 2 \cdot 3} \left(\frac{\sqrt{-1}}{16\sqrt{(3n^3)}} \right)^3 + \dots \right\}.$$
 (14)

Here, U denotes a solution to the Airy equation, $\frac{d^2U}{dn^2} + \frac{n}{3}U = 0$.

$$U = An^{-\frac{1}{4}} e^{\frac{2}{3}\sqrt{-\frac{n^3}{3}}} \left\{ 1 - \frac{1 \cdot 5}{1} \frac{\sqrt{-1}}{16\sqrt{(3n^3)}} + \frac{1 \cdot 5 \cdot 7 \cdot 11}{1 \cdot 2} \left(\frac{\sqrt{-1}}{16\sqrt{(3n^3)}} \right)^2 - \frac{1 \cdot 5 \cdot 7 \cdot 11 \cdot 13 \cdot 17}{1 \cdot 2 \cdot 3} \left(\frac{\sqrt{-1}}{16\sqrt{(3n^3)}} \right)^3 + \dots \right\}.$$
 (14)

Secondly, suppose *n* negative, and equal to -n'. Then, writing -n' for *n* in (14), and changing the arbitrary constant, and the sign of the radical, we get

$$U = Cn'^{-\frac{1}{4}} e^{-\frac{2}{3}\sqrt{n'^{2}}} \left\{ 1 - \frac{1 \cdot 5}{1 \cdot 16 (3n^{3})^{\frac{1}{2}}} + \frac{1 \cdot 5 \cdot 7 \cdot 11}{1 \cdot 2 \cdot 16 \cdot 3n^{3}} - \dots \right\}, \quad (17)$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ ○○○

Stokes Behavior of the Airy Function

$$\operatorname{Ai}(z) = \frac{1}{2\pi} \int_{\gamma} e^{i(zt - \frac{1}{3}t^3)} dt, \quad \operatorname{Im}(\gamma) = (-\infty e^{-\frac{2}{3}i\pi}, 0] \cup [0, \infty)$$

(ロ)、

Stokes Behavior of the Airy Function

$$\operatorname{Ai}(z) = \frac{1}{2\pi} \int_{\gamma} e^{i(zt - \frac{1}{3}t^3)} dt, \quad \operatorname{Im}(\gamma) = (-\infty e^{-\frac{2}{3}i\pi}, 0] \cup [0, \infty)$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三回 のへぐ

The Stokes behavior of the Airy function:



Stokes Behavior of the Airy Function

$$\operatorname{Ai}(z) = \frac{1}{2\pi} \int_{\gamma} e^{i(zt - \frac{1}{3}t^3)} dt, \quad \operatorname{Im}(\gamma) = (-\infty e^{-\frac{2}{3}i\pi}, 0] \cup [0, \infty)$$

The Stokes behavior of the Airy function:



▲□▶ ▲□▶ ▲□▶ ▲□▶ □ □ のへで

Airy Function Asymptotics

Asymptotic Expansions:

$$\operatorname{Ai}(z) \sim \frac{1}{2\sqrt{\pi}} z^{-\frac{1}{4}} e^{-\zeta} \sum_{n=0}^{\infty} (-1)^n \frac{c_n}{\zeta^n}$$

$$|\arg(z)| < \pi$$

Airy Function Asymptotics

Asymptotic Expansions:

$$\begin{split} \operatorname{Ai}(z) &\sim \frac{1}{2\sqrt{\pi}} z^{-\frac{1}{4}} e^{-\zeta} \sum_{n=0}^{\infty} (-1)^n \frac{c_n}{\zeta^n} & |\operatorname{arg}(z)| < \pi \\ \operatorname{Ai}(-z) &\sim \frac{1}{\sqrt{\pi}} z^{-\frac{1}{4}} \left(\sin(\zeta + \frac{\pi}{4}) \sum_{n=0}^{\infty} (-1)^n \frac{c_{2n}}{\zeta^{2n}} \\ &- \cos(\zeta + \frac{\pi}{4}) \sum_{n=0}^{\infty} (-1)^n \frac{c_{2n+1}}{\zeta^{2n+1}} \right) & |\operatorname{arg}(z)| < \frac{2}{3} \pi \end{split}$$

Airy Function Asymptotics

Asymptotic Expansions:

$$\begin{split} \operatorname{Ai}(z) &\sim \frac{1}{2\sqrt{\pi}} z^{-\frac{1}{4}} e^{-\zeta} \sum_{n=0}^{\infty} (-1)^n \frac{c_n}{\zeta^n} & |\operatorname{arg}(z)| < \pi \\ \operatorname{Ai}(-z) &\sim \frac{1}{\sqrt{\pi}} z^{-\frac{1}{4}} \left(\sin(\zeta + \frac{\pi}{4}) \sum_{n=0}^{\infty} (-1)^n \frac{c_{2n}}{\zeta^{2n}} \\ &- \cos(\zeta + \frac{\pi}{4}) \sum_{n=0}^{\infty} (-1)^n \frac{c_{2n+1}}{\zeta^{2n+1}} \right) & |\operatorname{arg}(z)| < \frac{2}{3} \pi \end{split}$$

Notation:

$$\zeta = \frac{2}{3}z^{\frac{3}{2}}, \quad c_0 = 1, \quad c_n = \frac{\Gamma(3n + \frac{1}{2})}{54^n n! \Gamma(n + \frac{1}{2})} = \frac{(2n+1)(2n+3)...(6n-1)}{216^n n!}$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 のへぐ

Asymptotics on the Real Line

Ai(x) ~
$$\begin{cases} \frac{1}{2\sqrt{\pi}} x^{-\frac{1}{4}} e^{-\frac{2}{3}x^{\frac{3}{2}}} & x > 0\\ \frac{1}{\sqrt{\pi}} (-x)^{-\frac{1}{4}} \sin\left(\frac{2}{3}(-x)^{\frac{3}{2}} + \frac{\pi}{4}\right) & x < 0 \end{cases}$$

(ロ)、(型)、(E)、(E)、(E)、(O)へ(C)

Asymptotics on the Real Line

$$\operatorname{Ai}(x) \sim \begin{cases} \frac{1}{2\sqrt{\pi}} x^{-\frac{1}{4}} e^{-\frac{2}{3}x^{\frac{3}{2}}} & x > 0\\ \frac{1}{\sqrt{\pi}} (-x)^{-\frac{1}{4}} \sin\left(\frac{2}{3}(-x)^{\frac{3}{2}} + \frac{\pi}{4}\right) & x < 0 \end{cases}$$



◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 のへで

Airy Function in the Complex Plane

Complex plots of the approximations and where they agree.

 $\operatorname{Ai}(z)$

(Entire)



▲□▶ ▲圖▶ ▲臣▶ ▲臣▶ 三臣 - のへで

Airy Function in the Complex Plane

Complex plots of the approximations and where they agree.



▲□▶ ▲□▶ ▲□▶ ▲□▶ □ ● のへで

Complex plots of the approximations and where they agree.



▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

Example Two: WKB Solution to a Helmholtz Equation

Consider this one-dimensional Helmholtz Equation:

$$\frac{d^2u}{dz^2}(z) = k^2 R^2(z)u(z)$$

◆□▶ ◆□▶ ◆∃▶ ◆∃▶ ∃ ∽のへで

Consider this one-dimensional Helmholtz Equation:

$$\frac{d^2u}{dz^2}(z) = k^2 R^2(z)u(z)$$

Medium of Varying Refractive Index μ :

$$R(z) = i\mu(z), \quad \mu(x) > 0, \quad \mu(x) \to 1 \text{ as } x \to \infty$$

Consider this one-dimensional Helmholtz Equation:

$$\frac{d^2u}{dz^2}(z) = k^2 R^2(z)u(z)$$

Medium of Varying Refractive Index μ :

$$R(z) = i\mu(z), \quad \mu(x) > 0, \quad \mu(x) \to 1 \text{ as } x \to \infty$$

Exponentially weak reflections arise:

M. Berry

We are able to answer this question [i.e. where and how does the reflected wave arise on the x-axis], because the birth of a reflection is simply the switching-on of a subdominant multiplier. Stokes lines arise from zeroes of R, say z_j , wherein:

$$\operatorname{Im} w(z) = 0, \quad w(z) := \int_{z_j}^z R(t) dt$$

◆□▶ ◆□▶ ◆∃▶ ◆∃▶ ∃ ∽のへで

Stokes lines arise from zeroes of R, say z_j , wherein:

$$\operatorname{Im} w(z) = 0, \quad w(z) := \int_{z_j}^z R(t) dt$$

M. Berry

Stokes lines lie at the heart of the asymptotics of [this equation.] They are the locus of greatest disparity between the dominant and subdominant fundamental phase-integral approximate solutions attached to z_i :

$$u_{\pm} \approx \exp(\pm kw(z))/R^{\frac{1}{2}}(z).$$

うして ふゆ く は く は く む く し く

Dominant u_+ corresponds to the incident wave, Subdominant u_- to the reflected wave.

$$u_{\pm} \approx \exp(\pm kw(z))/R^{\frac{1}{2}}(z).$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 のへぐ

Dominant u_+ corresponds to the incident wave, Subdominant u_- to the reflected wave.

$$u_{\pm} \approx \exp(\pm kw(z))/R^{\frac{1}{2}}(z).$$

WKB Approximate Solution:

$$u(z) \approx a_+(z)u_+(z) + a_-(z)u_-(z)$$

Across a Stokes line, the multiplier a_{-} jumps by ia_{+} .

Suppose a quantum particle moves freely in planar region \mathscr{B} with reflection at the boundary $\partial \mathscr{B}$.

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 のへぐ

Suppose a quantum particle moves freely in planar region \mathscr{B} with reflection at the boundary $\partial \mathscr{B}$.

Consider the Dirichlet eigenvalue problem:

$$\begin{cases} -\Delta \phi_n(r) = E_n \phi_n(r), & r = (x, y) \in \mathscr{B} \\ \phi_n = 0 & r \in \partial \mathscr{B} \end{cases}$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 のへぐ

Suppose a quantum particle moves freely in planar region \mathscr{B} with reflection at the boundary $\partial \mathscr{B}$.

Consider the Dirichlet eigenvalue problem:

$$\begin{cases} -\Delta \phi_n(r) = E_n \phi_n(r), & r = (x, y) \in \mathscr{B} \\ \phi_n = 0 & r \in \partial \mathscr{B} \end{cases}$$

Regularized Resolvent:

$$g(s) = \lim_{N \to \infty} \left(\sum_{n=1}^{N} \frac{1}{E_n + s^2} - \frac{\mathscr{A}}{4\pi} \log\left(\frac{E_N}{s^2}\right) \right)$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ □ のQ@

Example Three: Quantum Billiards

Weyl Expansion:

$$g(s) = \sum_{r=1}^{\infty} \frac{c_r}{s^r}$$

Example Three: Quantum Billiards

Weyl Expansion:

$$g(s) = \sum_{r=1}^{\infty} \frac{c_r}{s^r}$$

One may truncate to the least term, obtaining an exponentially small remainder. $$\sp{*}$$

$$g(s) = \sum_{r=1}^{r^*} \frac{c_r}{s^r} + R(s)$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

Example Three: Quantum Billiards

Weyl Expansion:

$$g(s) = \sum_{r=1}^{\infty} \frac{c_r}{s^r}$$

One may truncate to the least term, obtaining an exponentially small remainder.

$$g(s) = \sum_{r=1}^{r^*} \frac{c_r}{s^r} + R(s)$$

Across Stokes lines, the remainder changes behavior becoming oscillatory, then large.

$$e^{-s} \longrightarrow e^{-is} \longrightarrow e^{s}, s \in \mathbb{R}^{+}$$

うしゃ ふゆ きょう きょう うくの

(ロ)、

Quantum Billiards Examples

Quantum Billiards Examples

■ Spectral Resurgence



- Quantum Billiards Examples
- Spectral Resurgence
- Deducing Stokes behavior from expansions

- Quantum Billiards Examples
- Spectral Resurgence
- Deducing Stokes behavior from expansions

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三回 のへぐ

And more!

- G. B. Airy, "On the Intensity of Light in the neighbourhood of a Caustic," Trans. Cambridge Phil. Soc. Vol. 6, Pt. 3, 397-402 (1838)
- G. G. Stokes, "On the numerical Calculation of a Class of Definite Integrals and Infinite Series," Trans. Cambridge Phil. Soc. Vol. 9 Pt. 2, 166-187 (1850)
- G. G. Stokes, "On the discontinuity of arbitrary constants that appear as multipliers of semi-convergent series (A letter to the Editor)," Acta Math. Stockholm 26, 393-397 (1902)

References

- A. B. Olde Daalhuis, S. J. Chapman, J. R. King, J. R. Ockendon, and R. H. Tew, "Stokes phenomenon and matched asymptotic expansions," SIAM J. Appl. Math, Vol. 55, No. 6, 1469-1483 (1995)
- M. V. Berry, "Uniform asymptotic smoothing of Stokes' discontinuities," Proc. R. Soc. Lond. A 422, 7-21 (1989)
- M. V. Berry, "Waves near Stokes lines," Proc. R. Soc. Lond. A 427, 265-280 (1990)
- M. V. Berry and C. J. Howls, "High orders of the Weyl Expansion for quantum billiards: resurgence of periodic orbits, and the Stokes phenomenon," Proc. R. Soc. Lond. A 447, 527-555 (1994)
- O. Costin, Asymptotics and Borel Summability, Chapman & Hall/CRC Press, 2009.
- Abramowitz, M. and Stegun, I. A., eds. Handbook of Mathematical Functions with Formulas, Graphs, and Mathematical Tables, "Chapter 10," Applied Mathematics Series, 55, (Ninth reprint, 1983). Washington D.C.; New York: United States Department of Commerce, National Bureau of Standards; Dover Publications. p. 446-447

- Mika-Pekka Markkanen, Supernumerary Rainbows [Photograph], 23 May 2010, Wikimedia Commons. License: CC BY-SA 4.0.
- Les Cowley, Bow from 0.75mm diameter drops illuminated by a distant point source [Simulated Image, Cropped], n.d., Atmospheric Optics, atopics.co.uk. Copyright of Les Cowley, reproduced under Fair Use clause.
- Geek3, Airy Ai Asymptotic [SVG, converted to PNG], 7
 Feb. 2015, Wikimedia Commons. License: CC BY 3.0.