From Rainbows to Resurgence: Asymptotics of the Airy Function

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Supernumerary, or Spurious, Rainbow



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- In 1841, W. H. Miller observed/measured 30 dark bands for the primary bow
- In 1850, Stokes employed another method, managing to calculate 50 zeroes!

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$$u = e^{-i\frac{\pi}{6}} \int_0^\infty e^{-(z^3 - pz)} dx = U - iU'$$

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As noted by Stokes, though, this is "not convenient" when n becomes large.

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$$U = e^{\frac{2}{3}\sqrt{-\frac{n^{3}}{3}}} \left(An^{\alpha} + Bn^{\beta} + Cn^{\gamma} + ...\right)$$

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From the equation, he deduced that:

$$U = An^{-\frac{1}{4}}e^{\frac{2}{3}\sqrt{-\frac{n^3}{3}}} \left(1 - \frac{1 \cdot 5 \cdot i}{16\sqrt{3n^2}} + \frac{1 \cdot 5 \cdot 7 \cdot 11}{1 \cdot 2} \left(\frac{i}{16\sqrt{3n^2}}\right)^2 + \dots\right)$$

and similarly for negative values n' = -n.

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Just one small problem... the series diverges.

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Here's what Stokes' original series looks like:

$$U = An^{-\frac{1}{4}} e^{\frac{2}{3}\sqrt{-\frac{n^{3}}{3}}} \left\{ 1 - \frac{1.5}{1} \frac{\sqrt{-1}}{16\sqrt{(3n^{3})}} + \frac{1.5.7.11}{1.2} \left(\frac{\sqrt{-1}}{16\sqrt{(3n^{3})}} \right)^{2} - \frac{1.5.7.11.13.17}{1.2.3} \left(\frac{\sqrt{-1}}{16\sqrt{(3n^{3})}} \right)^{3} + \dots \right\}.$$
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 (14)

Secondly, suppose *n* negative, and equal to -n'. Then, writing -n' for *n* in (14), and changing the arbitrary constant, and the sign of the radical, we get

$$U = Cn'^{-\frac{1}{4}} e^{-\frac{2}{3}\sqrt{\frac{n'^2}{3}}} \left\{ 1 - \frac{1 \cdot 5}{1 \cdot 16 (3n^3)^{\frac{1}{2}}} + \frac{1 \cdot 5 \cdot 7 \cdot 11}{1 \cdot 2 \cdot 16 \cdot 3n^3} - \dots \right\}, \quad (17)$$

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When n or n' is at all large, the series [for U] are at first rapidly convergent, but they are ultimately in all cases hypergeometrically divergent. Notwithstanding this divergence, we may employ the series in numerical calculation, provided we do not take in the divergent terms.

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This is exactly what in modern days we would deem optimal (or least) truncation of an asymptotic expansion.

The expansion for U is different for negative n (what we now call the first example of the Stokes phenomenon.)

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This mode of determining the constant [in front of the negative U expansion] is anything but satisfactory. I have endeavored in vain to deduce the leading term in U for n negative from the integral itself...

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Today, we do know how to deduce the behavior for all n from the integral itself!

The Airy Function (Real Inputs)



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The method of steepest descent for an integral of the form:

$$\int_C f(t)e^{xg(t)}dt, \quad x \gg 1$$

involves deforming the integration contour to pass along the direction of steepest descent (i.e. parallel to $-\nabla u$, where g = u + iv) to pass by saddle points (viz. near where the integral is maximal, or at least less rapid oscillations cause less cancellation.)

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(In particular, this enables one to use Laplace's method to estimate the integral.)

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Thus, we write:

$$\operatorname{Ai}(z) = \frac{1}{2\pi i} \int_{\gamma} e^{\frac{t^3}{3} - zt} dt, \qquad \operatorname{Im}(\gamma) = (\infty e^{-\frac{\pi i}{3}}, 0] \cup [0, \infty e^{\frac{\pi i}{3}})$$

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Airy Function Asymptotics

Asymptotic Expansions:

$$\operatorname{Ai}(z) \sim \frac{1}{2\sqrt{\pi}} z^{-\frac{1}{4}} e^{-\zeta} \sum_{n=0}^{\infty} (-1)^n \frac{c_n}{\zeta^n}$$

$$|\arg(z)| < \pi$$

Airy Function Asymptotics

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$$\begin{split} \operatorname{Ai}(z) &\sim \frac{1}{2\sqrt{\pi}} z^{-\frac{1}{4}} e^{-\zeta} \sum_{n=0}^{\infty} (-1)^n \frac{c_n}{\zeta^n} & |\operatorname{arg}(z)| < \pi \\ \operatorname{Ai}(-z) &\sim \frac{1}{\sqrt{\pi}} z^{-\frac{1}{4}} \left(\sin(\zeta + \frac{\pi}{4}) \sum_{n=0}^{\infty} (-1)^n \frac{c_{2n}}{\zeta^{2n}} \\ &- \cos(\zeta + \frac{\pi}{4}) \sum_{n=0}^{\infty} (-1)^n \frac{c_{2n+1}}{\zeta^{2n+1}} \right) & |\operatorname{arg}(z)| < \frac{2}{3} \pi \end{split}$$

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Notation:

$$\zeta = \frac{2}{3}z^{\frac{3}{2}}, \quad c_0 = 1, \quad c_n = \frac{\Gamma(3n + \frac{1}{2})}{54^n\Gamma(n+1)\Gamma(n+\frac{1}{2})}$$

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Asymptotics on the Real Line

Ai(x) ~
$$\begin{cases} \frac{1}{2\sqrt{\pi}} x^{-\frac{1}{4}} e^{-\frac{2}{3}x^{\frac{3}{2}}} & x > 0\\ \frac{1}{\sqrt{\pi}} (-x)^{-\frac{1}{4}} \sin\left(\frac{2}{3}(-x)^{\frac{3}{2}} + \frac{\pi}{4}\right) & x < 0 \end{cases}$$

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Asymptotics in the Complex Plane

Complex plots of the approximations and where they agree.

 $\operatorname{Ai}(z)$

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The Airy function is governed by the asymptotic expansion:

$$\varphi_{\rm Ai}(z) = \sum_{n=0}^{\infty} \frac{a_n}{z^n} = \sum_{n=0}^{\infty} \left(-\frac{3}{4}\right)^n \frac{\Gamma(n+\frac{1}{6})\Gamma(n+\frac{5}{6})}{2\pi\Gamma(n+1)} \frac{1}{z^n}$$

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Resumming the Asymptotic Series

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(This is the same formula as before, but with a change of variables to make growth akin to $\Gamma(n)$ more manifest.)

$$a_n = \left(-\frac{3}{4}\right)^n \frac{\Gamma(n+\frac{1}{6})\Gamma(n+\frac{5}{6})}{2\pi\Gamma(n+1)} = \left(-\frac{2}{3}\right)^{-n} \frac{\Gamma(3n+\frac{1}{2})}{54^n\Gamma(n+1)\Gamma(n+\frac{1}{2})}$$

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More remarks:

• φ_{Ai} is factorially divergent (of Gevrey-class one.)

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 z = k^{3/2}/2 is a natural change of variables for ensuing resummation.

Interlude: Borel Summation Schematic



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Borel Summation Example



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Airy Series: Borel Summation

• The minor of φ_{Ai} is its (formal) Borel transform, forgetting the constant term:

$$\tilde{\varphi}_{\mathrm{Ai}} := \mathcal{B}[\varphi_{\mathrm{Ai}}] = \sum_{n=1}^{\infty} a_n \frac{\zeta^{n-1}}{(n-1)!}$$

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- $\tilde{\varphi}_{Ai}$ extends analytically to the universal cover of $\mathbb{C} \setminus \{0, -\frac{4}{3}\}$
- For any direction θ not along the negative real axis, the following converges for $\operatorname{Re}(ze^{i\theta}) > 0$:

$$S_{\theta}\varphi_{\mathrm{Ai}}(z) := a_0 + \mathcal{L}_{\theta}\mathcal{B}[\varphi_{\mathrm{Ai}}](z) = a_0 + \int_{0}^{\infty e^{i\theta}} \tilde{\varphi}_{\mathrm{Ai}}(\zeta) e^{-z\zeta} d\zeta$$

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Where before:

$$\operatorname{Ai}(k) \sim \frac{1}{2\sqrt{\pi}} k^{-\frac{1}{4}} e^{-\frac{2}{3}k^{\frac{3}{2}}} \varphi_{\operatorname{Ai}}(k^{\frac{3}{2}})$$

We now have:

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One can rotate the direction of summation for new regions of validity.

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Rotation of summation is fine up until one encounters the singularity on the negative real line.

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FIGURE 2. Right and left Borel-resummation.

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One can compares right and left-resummations, since

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$$S_{-\pi^{-}}\varphi_{Ai}(z) = S_{-\pi^{+}}\varphi_{Ai}(z) + \int_{\gamma} \widetilde{\varphi_{Ai}}(\zeta) e^{-z\zeta} d\zeta$$

The Hankel contour γ can be expressed using the so-called alien derivative:

$$\int_{\gamma} \tilde{\varphi}_{\mathrm{Ai}}(\zeta) e^{-z\zeta} d\zeta = e^{+\frac{4}{3}z} S_{-\pi} \left(\Delta_{-\frac{4}{3}}^{z} \varphi_{\mathrm{Ai}} \right) (z)$$

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In this case,

$$\Delta_{-\frac{4}{3}}^{z}\varphi_{\mathrm{Ai}} = -i\varphi_{\mathrm{Bi}}, \quad \varphi_{\mathrm{Bi}}(z) := \sum_{n=0}^{\infty} (-1)^{n} \frac{a_{n}}{z^{n}}$$

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Airy expansion when $|\arg(k) - \pi| < \frac{\pi}{3}, z = k^{\frac{3}{2}}$:

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$$\operatorname{Ai}(k) = \frac{1}{2\sqrt{\pi}} k^{-\frac{1}{4}} \left(e^{-\frac{2}{3}z} S_{-\frac{3\pi}{2}} \varphi_{\operatorname{Ai}}(z) + i e^{+\frac{2}{3}z} S_{-\frac{3\pi}{2}} \varphi_{\operatorname{Bi}}(z) \right)$$

Note the new exponential term that appeared.

Once can rewrite the LHS as the resummed version of the second expansion we saw previously.

Zeroes "Resurge" from the Original Expansion

• This expression can be employed to calculate the zeroes on the Airy function (viz. the Delabaere reference.)

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- This expression can be employed to calculate the zeroes on the Airy function (viz. the Delabaere reference.)
- This procedure follows the original calculations of Stokes, but now by manipulating the full series itself. The mathematical justification for the resummation comes from resurgent analysis.

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- This procedure follows the original calculations of Stokes, but now by manipulating the full series itself. The mathematical justification for the resummation comes from resurgent analysis.
- In particular, the behavior on the negative real line is manifestly contained in the expansion on the positive real line— an example of resurgence.

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